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# CONCERNING THE DISCONTINUOUS SOLUTION IN THE PROBLEM OF THE MINIMUM SURFACE OF REVOLUTION

BY HARRIS F. MACNEISH

1. In the problem of minimizing the integral

$$J = \int_{t_0}^{t_1} y \sqrt{x'^2 + y'^2} dt,$$

where the admissible curves are all "ordinary"\* curves which can be drawn in the upper half plane ( $y \geq 0$ ) from the given point  $A$  to the given point  $B$ , Euler's differential equation has two solutions:

(1) the catenaries

$$x = t ; \quad y = m \operatorname{ch} \frac{t - x_0'}{m} ; \quad (1)$$

(2) the straight lines

$$x = a ; \quad y = t. \quad (2)$$

The latter solution leads to the well-known "discontinuous solution" of the problem, first noticed by Goldschmidt,† which consists (see Figure 3) of

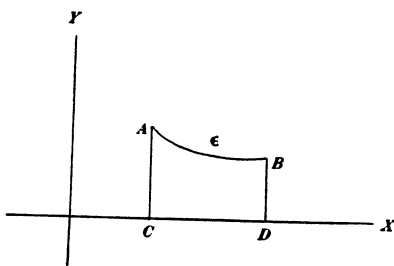


FIG. 3.

the perpendicular  $AC$  to the  $x$ -axis, the segment  $CD$  of the  $x$ -axis, and the perpendicular  $DB$ .

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\* In the terminology of Bolza, *Lectures on the Calculus of Variations*, p. 117.

† *Göttingen Prize Essay*, 1831.

The segment  $CD$  is not itself an extremal but it satisfies the boundary conditions for a minimum arising from the first variation (see Bolza, *Lectures on the Calculus of Variations*, p. 153).

Todhunter\* has given a simple geometrical sufficiency proof that  $I_{ACDB} < I_{AB}$  for every admissible curve  $AB = \mathcal{C}$  whose length is greater than  $|AC| + |DB|$  (see Figure 3).

We now consider a catenary of the system (1) passing through the point  $A (x_0, y_0)$ , namely :

$$y = m \operatorname{ch} \frac{x - x'_0}{m}, \quad (3)$$

where

$$y_0 = m \operatorname{ch} \frac{x_0 - x'_0}{m}; \quad (4)$$

and take on this catenary the point  $B (x_1, y_1)$ .

We propose to compare the areas of the surfaces of revolution generated by the arc  $AB$  of the catenary on the one hand and by the discontinuous solution  $ACDB$  on the other hand.

The value of the latter is

$$S_D = \pi (y_0^2 + y_1^2) = \pi m^2 \left\{ \operatorname{ch}^2 \frac{x_0 - x'_0}{m} + \operatorname{ch}^2 \frac{x_1 - x'_0}{m} \right\},$$

and the value of the former

$$S_C = \frac{\pi m}{2} \left\{ 2(x_1 - x_0) + m \operatorname{sh} \frac{2(x_1 - x'_0)}{m} - m \operatorname{sh} \frac{2(x_0 - x'_0)}{m} \right\}.$$

If we introduce the abbreviations

$$u_0 = \frac{x_0 - x'_0}{m}, \quad u_1 = \frac{x_1 - x'_0}{m} \quad (5)$$

we obtain for the difference

$$\begin{aligned} S_D - S_C &= \pi m^2 \{ (2u_0 + \operatorname{sh} 2u_0 + 2 \operatorname{ch}^2 u_0) - (2u_1 + \operatorname{sh} 2u_1 - 2 \operatorname{ch}^2 u_1) \} \\ &= \pi m^2 \{ (2u_0 + 1 + e^{2u_0}) - (2u_1 - 1 - e^{-2u_1}) \}. \end{aligned}$$

As the point  $B$  moves along the catenary in the direction of the increasing

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\* Todhunter, *Researches on the Calculus of Variations*, §§64, 65

$x$ , the quantities  $m$ ,  $x_0$ ,  $x'_0$ , and therefore also  $u_0$ , remain constant, while  $u_1$  increases. We therefore consider the function

$$\phi(u_1) = 2u_1 - 1 - e^{-2u_1},$$

whence

$$\phi'(u_1) = 2(1 + e^{-2u_1}) > 0.$$

Then  $\phi(u_1)$  increases with  $u_1$  and therefore  $S_D - S_C$  decreases with  $u_1$ . When  $B$  coincides with  $A$ ,  $u_1 = u_0$  and  $S_D - S_C = 4\pi m^2 \text{ch}^2 u_0 > 0$ ; as  $x_1$  and therefore  $u_1$  approaches  $+\infty$ ,  $S_D - S_C$  approaches  $-\infty$ . Hence there exists one and but one value of  $u_1$  for which  $S_D - S_C = 0$ ; that is:

*There exists one and but one point  $A''$  on the catenary for which the discontinuous solution has the same value as the catenary solution, i. e. for which*

$$2u_0 + 1 + e^{2u_0} = 2u_1 - 1 - e^{-2u_1}. \quad (6)$$

If  $B$  lies between  $A$  and  $A''$ , the catenary solution has a smaller value than the discontinuous solution, while if  $B$  lies beyond  $A''$ , the discontinuous solution has the smaller value.

The point  $A''$  always lies between  $A$  and its conjugate  $A'$ ; for by Lindelöf's theorem\* at the point  $A'$  (see Figure 4)

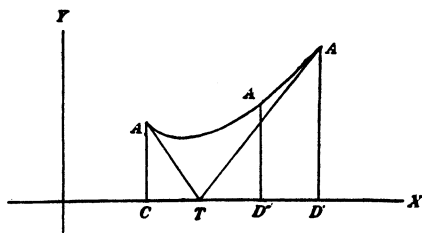


FIG. 4.

the area  $(AA')^\dagger$  is equal to the area generated by the two tangents  $AT$  and  $A'T$  which moreover intersect on the  $x$ -axis. But the area  $(ACD'A')$  is less than the area  $(ATA')$ , and therefore also less than the area  $(AA')$ . Hence according to the above result the point  $A'$  must lie beyond  $A''$ .

**2.** If we construct the point  $A''$  for every catenary of the set (1) through the point  $A$ , the points  $A''$  describe a curve which we propose to study in this section.

\* Hancock, "On the number of catenaries that may be drawn through two given points," *ANNALS OF MATHEMATICS*, ser. 1, vol 10, p. 159, §16; also *Calculus of Variations*, chapter III.

† I. e., the area generated by the arc  $AA'$ .

The coordinates of the point  $A''$  for a given catenary are determined by the equations

$$x_1 = x_0 + y_0 \frac{u_1 - u_0}{\text{ch } u_0}, \quad (7)$$

$$y_1 = y_0 \frac{\text{ch } u_1}{\text{ch } u_0}, \quad (8)$$

when  $u_0$  and  $u_1$  are connected by the relation

$$2u_0 + 1 + e^{2u_0} = 2u_1 - 1 - e^{-2u_1}. \quad (9)$$

Hence we obtain the required locus if we eliminate  $u_0$  from the equations (7), (8), (9) or else express  $u_0, u_1$  in terms of a variable  $t$ . It is convenient to choose for  $t$  the common value of the two sides of equation (9) :

$$t = 2u_0 + 1 + e^{2u_0} = 2u_1 - 1 - e^{-2u_1}. \quad (10)$$

We shall now consider the two curves in the  $u$ - $t$  plane represented by equation (10). In the first place,

$$\frac{dt}{du_0} = 2(1 + e^{2u_0}) > 0;$$

therefore  $t$  increases with  $u_0$ . Secondly, (11)

$$\frac{d^2t}{du_0^2} = 4e^{2u_0} > 0;$$

therefore the curve is convex to the  $u_0$ -axis. Again, (12)

$$\frac{dt}{du_1} = 2(1 + e^{-2u_1}) > 0;$$

therefore  $t$  increases with  $u_1$ . Finally, (13)

$$\frac{d^2t}{du_1^2} = -4e^{-2u_1} < 0;$$

therefore the curve is concave to the  $u_1$ -axis. (14)

Since both functions increase with  $u_0, u_1$  respectively, it follows that  $u_0$  and  $u_1$  are single valued functions of  $t$ . In order to obtain corresponding values of  $u_0$  and  $u_1$  we construct the two graphs upon the same set of axes from the following table :

TABLE I

$u_0$	$t$	$u_1$	$t$
.00	2.00	.00	-2.00
.20	2.89	.20	-1.27
.40	4.03	.40	-0.65
.60	5.52	.60	-0.10
.80	7.55	.80	+0.40
1.00	10.39	1.00	+0.86
:	:	1.50	1.95
$\infty$	$\infty$	2.00	2.98
-.20	1.27	2.50	3.99
-.40	0.65	3.00	4.99
-.60	0.10	asymptotic to line	
-.80	-0.40	$t = 2u_1 - 1$	
-1.00	-0.86	for positive	values of $u_1$
-1.50	-1.95	-.20	-2.89
-2.00	-2.98	-.40	-4.03
asymptotic to straight		-0.60	-5.52
line $t = 2u_0 + 1$		-0.80	-7.55
for negative values of $u_0$		-1.00	-10.39
		:	:
		$-\infty$	$-\infty$

The left branch is the curve  $t = 2u_0 + 1 + e^{2u_0}$ , while the right branch is the curve  $t = 2u_1 - 1 - e^{-u_1}$ . (See Figure 5.)

We next compute the derivatives  $dx_1/dt$  and  $dy_1/dt$  by means of equations (7), (8), (10), (12), and (14). First,

$$\frac{dx_1}{dt} = y_0 \frac{\text{ch } u_0 \left( \frac{du_1}{dt} - \frac{du_0}{dt} \right) - (u_1 - u_0) \text{sh } u_0 \frac{du_0}{dt}}{\text{ch}^2 u_0} \quad (15)$$

$$= y_0 \frac{\text{ch}^2 u_0 \text{sh } u_1 (\text{ch } u_1 + \text{sh } u_1) - \text{ch}^2 u_1 \text{sh } u_0 (\text{ch } u_0 - \text{sh } u_0)}{4 \text{ch } u_1 \text{ch}^3 u_0 (\text{ch } u_1 + \text{sh } u_1)} \quad (16)$$

$$= y_0 \frac{\text{ch}^2 u_0 \text{sh}^2 u_1 + \text{ch}^2 u_1 \text{sh}^2 u_0 + \text{ch } u_1 \text{ch } u_0 \text{sh}(u_1 - u_0)}{4 \text{ch } u_1 \text{ch}^3 u_0 (\text{ch } u_1 + \text{sh } u_1)} \quad (17)$$

Here the hyperbolic cosine is always positive,  $y_0$  is positive, and

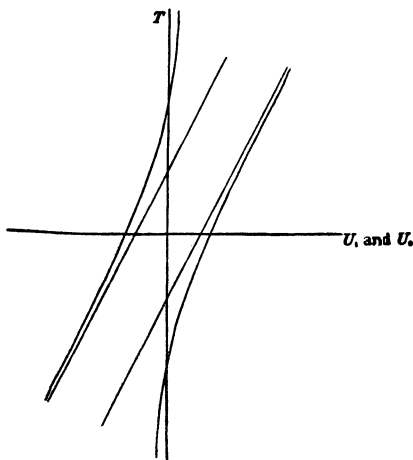


FIG. 5.

$u_1 - u_0 = \frac{1}{2} (2 + e^{2u_0} + e^{-2u_1}) > 0$  from (9); therefore  $dx_1/dt$  is positive, and  $t$  increases with  $x_1$ . Again,

$$\frac{dy_1}{dt} = y_0 \frac{\text{ch } u_0 \text{sh } u_1 \frac{du_1}{dt} - \text{ch } u_1 \text{sh } u_0 \frac{du_0}{dt}}{\text{ch}^2 u_0} \quad (18)$$

$$= y_0 \frac{\text{ch}^2 u_0 \text{sh } u_1 (\text{ch } u_1 + \text{sh } u_1) - \text{ch}^2 u_1 \text{sh } u_0 (\text{ch } u_0 - \text{sh } u_0)}{4 \text{ch } u_1 \text{ch}^3 u_0} \quad (19)$$

$$= y_0 \frac{\text{ch}^2 u_0 \text{sh}^2 u_1 + \text{ch}^2 u_1 \text{sh}^2 u_0 + \text{ch } u_1 \text{ch } u_0 \text{sh}(u_1 - u_0)}{4 \text{ch } u_1 \text{ch}^3 u_0}. \quad (20)$$

Therefore  $\frac{dy_1}{dt} > 0$ , and  $t$  increases with  $y_1$ . Then from (17) and (20),

$$\frac{dy_1}{dx_1} = \operatorname{ch} u_1 + \operatorname{sh} u_1 = e^{u_1} > 0, \quad (21)$$

and  $y_1$  increases with  $x_1$ . Further,

$$\frac{d_2 y_1}{dx_1^2} = e^{u_1} \frac{du_1}{dx_1} = e^{u_1} \cdot \frac{du_1}{dt} \cdot \frac{dt}{dx_1}, \quad (22)$$

where  $e^{u_1} > 0$ ,  $\frac{du_1}{dt} > 0$  from (13), and  $\frac{dt}{dx_1} > 0$  from (17).

Therefore  $\frac{d_2 y_1}{dx_1^2} > 0$  and the curve represented by equations (7), (8), (9)

is convex to the  $x$ -axis.

As  $t$  increases from  $-\infty$  to  $+\infty$ ,  $x_1$  and  $y_1$  both increase continually; as  $t$  approaches  $-\infty$ , i. e. as  $u_0$  and  $u_1$  (see table I) approach  $-\infty$ ,  $x_1$  and  $y_1$  approach zero. This is proved as follows:

$$\begin{aligned} x_1 &= y_0 \frac{u_1 - u_0}{\operatorname{ch} u_0} = y_0 \frac{2 + e^{2u_0} + e^{-2u_1}}{e^{u_0} + e^{-u_0}} \\ &= y_0 \frac{e^{-(2u_1 - u_0)} (2e^{2u_1} + e^{2u_0 + 2u_1} + 1)}{e^{2u_0} + 1}. \end{aligned} \quad (23)$$

Now  $u_0 - u_1 = -\frac{1}{2}(2 + e^{2u_0} + e^{-2u_1})$ , from (9); therefore

$$u_0 - 2u_1 = -\frac{1}{2}(2 + e^{2u_0} + 2u_1 + e^{-2u_1}).$$

Again,

$$2u_1 + e^{-2u_1} = \frac{1 + 2u_1 e^{2u_1}}{e^{2u_1}},$$

which approaches  $+\infty$  as  $u_1$  approaches  $-\infty$ , since  $2u_1 e^{2u_1}$  approaches 0. Therefore  $u_0 - 2u_1$  approaches  $-\infty$  as  $u_1$  and  $u_0$  approach  $-\infty$ . Then from equation (23),  $x_1$  approaches 0 as  $u_1$  and  $u_0$  approach  $-\infty$ .

$$\begin{aligned} \text{Moreover, } y_1 &= y_0 \frac{\operatorname{ch} u_1}{\operatorname{ch} u_0} = y_0 \frac{e^{u_1} + e^{-u_1}}{e^{u_0} + e^{-u_0}} \\ &= y_0 \frac{e^{-(u_1 - u_0)} (e^{2u_1} + 1)}{e^{2u_0} + 1}. \end{aligned}$$



Therefore  $y_1$  approaches 0 as  $u_1$  and  $u_0$  approach  $-\infty$ , since

$$u_1 - u_0 = \frac{1}{2} (2 + e^{2u_0} + e^{-2u_1})$$

approaches  $+\infty$ . Therefore the graph of our curve starts at the origin.

Similarly, as  $t$  approaches  $+\infty$ , i. e. as  $u_1$  and  $u_0$  approach  $+\infty$  (see table I),  $x_1$  and  $y_1$  approach  $+\infty$ ; for,

$$x_1 = y_0 \frac{u_1 - u_0}{\text{ch } u_0} = y_0 \frac{(2 + e^{2u_0} + e^{-2u_1})}{e^{u_0} + e^{-u_0}} = y_0 \frac{2e^{-u_0} + e^{u_0} + e^{-2u_1 - u_0}}{1 + e^{-2u_0}},$$

so that  $x_1$  approaches  $+\infty$  as  $u_1$  and  $u_0$  approach  $+\infty$ ;

and

$$y_1 = y_0 \frac{e^{u_1} + e^{-u_1}}{e^{u_0} + e^{-u_0}} = y_0 \frac{e^{u_1 - u_0} + e^{-u_1 - u_0}}{1 + e^{-2u_0}},$$

which approaches  $+\infty$  as  $u_1$  and  $u_0$  approach  $+\infty$ , since  $u_1 - u_0$  approaches  $+\infty$ .

Taking corresponding values of  $u_0$  and  $u_1$  from the graph given in Figure 5 we obtain the following table of corresponding values of  $x_1$  and  $y_1$ , for  $x_0 = 0$  and  $y_0 = 1$ .

TABLE II

$u_1$	$t$	$u_0$	$x_1$	$y_1$
$-\infty$	$-\infty$	$-\infty$	.00	.00
-1.00	-10.39	-5.14	.05	.02
.00	-2.00	-1.55	.63	.41
.50	-0.36	-0.80	.97	.84
1.00	+0.86	-0.33	1.26	1.46
1.50	1.88	-0.04	1.54	2.35
2.00	2.90	+0.21	1.75	3.68
2.50	3.91	0.39	1.96	5.69
3.00	4.93	0.55	2.12	8.66
.	.	.	.	.
.	.	.	.	.
$\infty$	$\infty$	$\infty$	$\infty$	$\infty$

With the aid of this table we construct the graph  $\mathcal{G}$  determined by equa-

tions (7), (8), (9), and on the same figure we draw the curve  $F$  given in Figure 2 of the preceding paper.

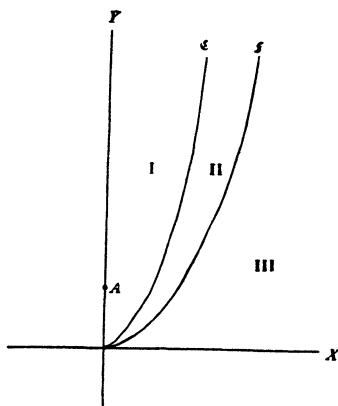


FIG. 6.

From our results in the two papers we know that from the point  $A$  to points of regions I and II two catenaries can be drawn with the  $x$ -axis as directrix; of these the upper one furnishes a relative minimum in the problem of the minimum surface of revolution;\* to points of curve  $F$  one catenary can be drawn, and to points of region III no catenaries can be drawn.

Then considering the two solutions of the problem of the minimum surface of revolution; i. e. the catenary solution and the discontinuous solution, we can conclude as follows:

- 1) *In the region I both the catenary solution and the discontinuous solution give a relative minimum, and the surface of the catenary is the smaller.*
- 2) *For points along the curve  $\mathcal{C}$  both solutions give a relative minimum and the surfaces are equal.*
- 3) *In the region II both solutions give a relative minimum and the discontinuous solution is the smaller.*
- 4) *Along the curve  $F$  the catenary solution does not furnish a minimum,† so that the discontinuous solution is the only solution.*
- 5) *In the region III the discontinuous solution is the only solution.*

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\* Todhunter, *Researches in the Calculus of Variations*, p. 57; Hancock, *Calculus of Variations*, chapters II and III.

† Lindelöf's Theorem; see Hancock, "On the number of catenaries which may be drawn through two fixed points," *ANNALS OF MATHEMATICS*, ser. 1, vol. 10, p. 159, or *Calculus of Variations*, chapter III.